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1-The multi-scale entropy analysis can unveil the self-similarity in time series.

2-The result of multi-scale entropy for fractional Gaussian noise is well modeled by a decreasing  $q$ -exponential function.

3-The Hurst exponent of a time series can be determined by the multi-scale entropy analysis.

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# Multi-scale entropy analysis and Hurst exponent

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## Abstract

Several methods exist for measuring the complexity in a system through analysis of its associated time series. Multi-scale entropy appears as a successful method on this matter. It has been applied in many disciplines with great achievements. For example, analysis of the bio-signals, we are able to diagnose various diseases. However, in most versions for the multi-scale entropy the examined time series is analyzed qualitatively. In this study, we try to present a quantitative picture for the multi-scale entropy analysis. Particularly, we focus on finding relation between the result of the multi-scale analysis and the Hurst exponent which quantifies the persistence in time series. For this purpose, the fractional Gaussian noise time series with different Hurst exponents are analyzed by the multi-scale entropy method and the results are fitted to a decreasing  $q$ -exponential function. We observe remarkable relation between the function parameters and Hurst exponent. This function can simulate the result of analysis for the white noise to the  $1/f$  noise.

*Keywords:* Multi-scale Entropy, Fractional Gaussian Noise, Hurst Exponent

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## 1. Introduction

Time series give many information about the examined system. Such information are mostly coded as self-similar patterns in time series. The ex-

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4 istence of these patterns is usually due to the long range interaction between  
5 system components and/or long term memory in their dynamics. Every pat-  
6 tern exhibits an order or regularity in the system and allows us to predict  
7 the system's behavior statistically. This is what most of the scientists called  
8 complexity [1]. Albeit, some other definitions of the complexity exist which  
9 differ with the above explanation. Most of them deal with the amount of  
10 information that we need for understanding of the system behavior. Re-  
11 garding to this perspective, the concept of entropy appears as a simple and  
12 powerful measure for quantifying the complexity. It can be found by simple  
13 googling many literature which discuss alternative methods for entropy esti-  
14 mation of the complexity in real world data. Here we try to reconcile these  
15 two treatments of the complexity.

16 Several methods have been proposed for identification of the self-similar  
17 pattern in time series [2]. The self-similarity may exist in time series graph  
18 which be revealed by the fractal and multi-fractal analysis [3, 4]. The scale-  
19 free distribution of values is another indication for self-similarity in time  
20 series [5]. The power law relationship for the auto-correlation function is  
21 the most well known feature which represents the self-similar pattern in time  
22 series. It deserves to note that the above aspects of self-similarity may be  
23 related to one another.

24 For the first time, It was H. E. Hurst a British hydrologist who pointed  
25 out to the self-similarity of a time series while studying on the optimum dam  
26 size for the Nile river in 1951. He developed a method for measuring the self-  
27 similarity which is known as rescale range analysis (RS) [6]. The result of RS  
28 analysis is expressed as a value between 0 and 1, which is called the Hurst  
29 exponent. There are many other methods that directly or indirectly compute  
30 the Hurst exponent for a time series. Now, we only name some of them and  
31 refer the interested readers to literature [2, 7]. As some instances, we briefly  
32 point out to the following methods, detrended fluctuation analysis (DFA)  
33 [8, 9, 3], power spectral analysis and its variants [10], wavelet method [11],  
34 methods based on the complex networks theory [12] (see also it's references).

35 In this work, we are intended to show that the multi-scale entropy (MSE)  
36 [13, 14, 15, 16, 17, 18, 19], can be also used for estimation of the Hurst  
37 exponent. For this end, we apply the MSE method to analyze the fractional  
38 Gaussian noise (FGN) time series with various Hurst exponents. Then we  
39 fit the obtained results by the decreasing  $q$ -exponential as a model function.  
40 We will find that the value of parameters are nicely related to the Hurst  
41 exponents.

42 We organize this paper as follows. In the next section we briefly review  
 43 the fractional Gaussian noise and meaning of the Hurst exponent. The basic  
 44 variant of multi-scale entropy is discussed in the third section. Fourth section  
 45 is devoted to the results and their interpretation. We summarize our work  
 46 at the final section.

## 47 2. Fractional Gaussian Noise

48 The fractional Brownian motion (FBM) is one of the most well known  
 49 stochastic processes which has been widely studied analytically [20]. It is  
 50 used in modelling various phenomena in science and engineering [21, 22].  
 51 Researchers are interested in studying and using FBM for its properties like  
 52 as self-similarity [21, 22].

53 The one dimensional FBM which we denote it as  $B_H(t)$ , is a non-  
 54 stationary stochastic process which starts at zero,  $B_H(0) = 0$ . The process  
 55 is known to have zero means,  $\langle B_H(t) \rangle = 0$ , and Gaussian distribution for its  
 56 increments. The auto-covariance function of the process is,

$$\langle B_H(t + \tau) B_H(t) \rangle = \frac{1}{2} (|t + \tau|^{2H} + |t|^{2H} - |\tau|^{2H}). \quad (1)$$

57 For simplicity, we assume that  $\langle B_H^2(1) \rangle = 1$ . The parameter  $H$  is the Hurst  
 58 exponent, a real value between 0 and 1 which determines the persistence of  
 59 process. For  $H = 0.5$  we have ordinary Brownian motion or Wiener process.

60 Dependence of auto-covariance function on  $t$  is the reason for non-stationary  
 61 feature of FBM. By using equation 1 in the case  $\tau = 0$ , we obtain,  $\langle$   
 62  $B_H^2(t) \rangle = t^{2H}$  or  $\langle B_H^2(at) \rangle = a^{2H} \langle B_H^2(t) \rangle$ . By simple calculation  
 63 it can be generalized to all other moments. This result indicates that the  
 64 FBM process is self-similar in distribution,  $B_H(at) \stackrel{d}{=} a^H B_H(t)$ .

65 As we mentioned earlier, FBM time series has non-stationary character  
 66 and is not suitable for modelling the stationary processes. The increments of  
 67 FBM as denoted as,  $\Delta B_H(t) = B_H(t + \Delta t) - B_H(t)$ , define another stochas-  
 68 tic process, named as fractional Gaussian noise (FGN). By using the auto-  
 69 covariance function of FBM and some algebraic manipulation, we can prove  
 70 that the auto-covariance of this new process is,

$$\langle \Delta B_H(t + \tau) \Delta B_H(t) \rangle = \frac{1}{2} (|\tau + \Delta t|^{2H} + |\tau - \Delta t|^{2H} - 2|\tau|^{2H}). \quad (2)$$

71 In the case,  $\tau \gg \Delta t$ , we can approximate the auto-covariance function as,

$$\langle \Delta B_H(t + \tau) \Delta B_H(t) \rangle \sim H(2H - 1) \tau^{H-2}. \quad (3)$$

72 The case  $H > 0.5$  demonstrates the existence of long range dependence in  
 73 series. The FGN also inherits the property of self-similarity from its parent;  
 74 by putting  $\tau = 0$  in equation 2 we arrive at the desired result for second  
 75 moment. The Gaussian distribution of  $\Delta B_H(t)$ , allows us to generalize the  
 76 obtained result easily to all other moments.

77 It is important to note that the Hurst exponent is an indicator for both  
 78 self-similarity and long range dependence in FGN time series.

### 79 3. Multi-scale Entropy

80 In statistical mechanics, entropy is a quantity for measuring disorder in a  
 81 system. A system may have many microscopic states and each state has  
 82 certain probability of occurrence. The goal of statistical mechanics is to  
 83 predict these probabilities. All the entropies like as Gibbs-Shannon, Renyi  
 84 and Tsallis, are expressed in terms of the state probabilities.

85 When disorder is increased in the system, this means that most of the  
 86 states are likely to occur then it is so hard to predict the state of system.  
 87 Entropy also increases in this case. If disorder decreases, some states will  
 88 be preferred and system becomes predictable, therefore entropy is decreased.  
 89 The spatial and temporal patterns are indication of the regularity or order  
 90 in a system. Any regularity makes the system predictable. Entropy can  
 91 measure the amount of predictability or in other word the complexity of  
 92 system.

93 The most challenging problem is the estimation of entropy for a given  
 94 system from its time series data. Here we only focus on the multi-scale  
 95 entropy analysis which is the most powerful method for investigating the  
 96 complexity of time series. In following we describe the method and refer the  
 97 interested readers for convincing statements to Refs. [13, 14].

98 Assume, we have a time series of length  $N$  which is denoted by series of  
 99 values,  $\{x_1, \dots, x_N\}$ .

100 The interaction between system and its environment may induce noises in  
 101 the system time series. Short range correlations in noises can be accumulated  
 102 and make a long range effect which is non-original. In the first step, it is  
 103 necessary to reduce the effect of unwanted noises and short range correlations

104 from signal by the coarse-graining procedure. In this step, we partitioned  
 105 time series into non-overlapping windows with equal length  $\tau$ , then the coarse  
 106 grained time series is constructed by averaging on data in each window,

$$x_i^{(\tau)} = \frac{1}{\tau} \sum_{j=1}^{\tau} x_{(i-1)\tau+j} \quad (4)$$

107 and putting them in a new sequence,  $\{x_1^{(\tau)}, \dots, x_{\lfloor n/\tau \rfloor}^{(\tau)}\}$ . In literature,  $\tau$  is  
 108 called the scale factor.

109 For the resulted time series we can define the  $m$ -dimensional vectors like  
 110 as,

$$\mathbf{X}_m^{(\tau)}(i) = \{x_i^{(\tau)}, \dots, x_{i+m-1}^{(\tau)}\}. \quad (5)$$

111 In the next step, we count the number of vector pairs that have distance  
 112 less than  $r$  and denote it by  $n_m(r, \tau)$ . We repeat the same computation for  
 113  $(m+1)$ -dimensional vector pairs and obtain  $n_{m+1}(r)$ . The sample entropy  
 114 is defined as,

$$S_E(m, r, \tau) = -\log(n_{m+1}(r, \tau)/n_m(r, \tau)). \quad (6)$$

115 It is clear that the sample entropy has zero or positive value, because  
 116  $n_{m+1}(r, \tau)$  is less than  $n_m(r, \tau)$ .

117 In final step we plot the sample entropy against the scale factor. For the  
 118 white noise we observe a decreasing behavior and it is proved analytically.  
 119 For the short range correlated signals it is expected to find the same behavior  
 120 as the white noise since coarse-graining procedure eliminates the short range  
 121 correlations. The long range dependence in time series does not change after  
 122 the coarse-graining operation then for  $1/f$  noise the sample entropy does not  
 123 vary with scale factor and remains constant. In the next section we analysis  
 124 the FGN data series with different Hurst exponent by MSE method.

## 125 4. Results

126 First of all, we should generate the needed FGN series with different  
 127 range of Hurst exponent for our analysis. The FGN samples are obtained  
 128 through differencing of the fractional Brownian motion data. There are sev-  
 129 eral ways for generating the Brownian motion time series with given Hurst  
 130 exponent. Among them, we choose three algorithms; Hosking [23], random  
 131 midpoint displacement [24] and Rambaldi-Pinazza [25]. These algorithms  
 132 employ different approaches for doing this. The Hosking algorithm uses the

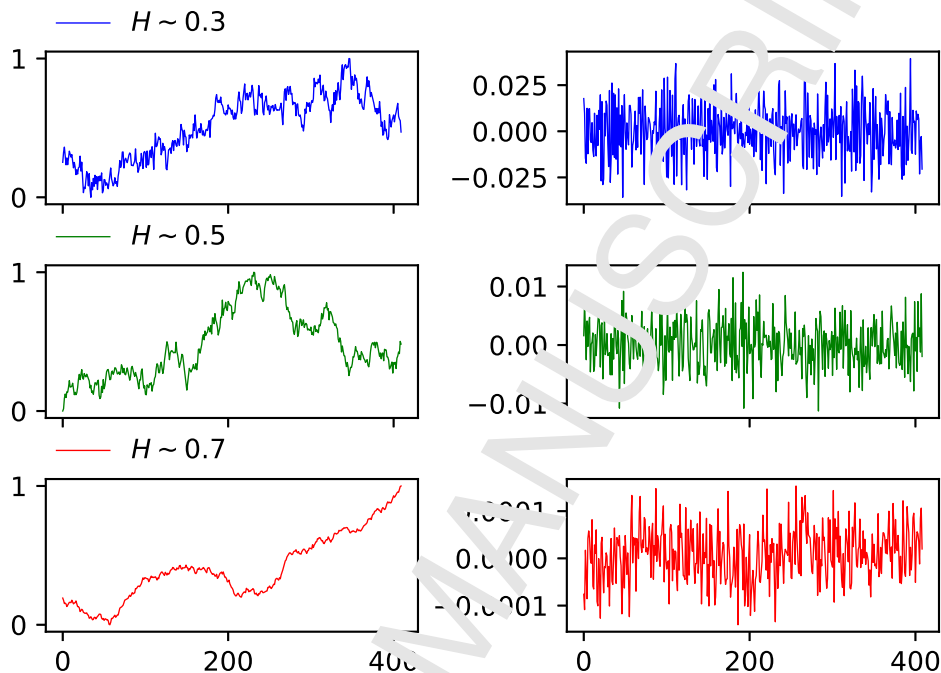


Figure 1: Part of the FBM (left) and FGN (right) time series for three values of the Hurst exponent,  $\sim 0.3$ ,  $\sim 0.5$  and  $\sim 0.7$ . For comparison purpose, the FBM time series are scaled to interval  $[0, 1]$ . Time series with small Hurst exponent is more denticulated.

133 FBM covariance property. In the random midpoint displacement method  
 134 the self similarity and Gaussian distribution of increments are used and finally the Rambaldi-Pinazza algorithm considers the integral representation  
 135 of FBM. Thus, we construct 48 samples with Hurst exponent between 0.15  
 136 and 0.9 and size of 65535. In construction procedure all data are scaled between 0 and 1. This scaling does not alter the Hurst exponent. Figure 1  
 137 shows the plot for FBM and its associated FGN series for three values of the  
 138 Hurst exponent,  $\sim 0.3$ ,  $\sim 0.5$  and  $\sim 0.7$ . We observe that by increasing the  
 139 Hurst exponent the difference between adjacent values in FBM time series  
 140 decreases. It should be noted that the procedure of generating FBM time series is not completely exact. Therefore it is necessary to measure the Hurst  
 141 exponent again for all samples data. Here we use DFA for this purpose.  
 142

143 Before analyzing the prepared FGN data by MSE, we should set two  
 144 parameters,  $m$  and  $r$ . The first one determines the dimension of vectors, see  
 145 equation 5, and the latter defines a threshold for closeness of vectors. Here  
 146  
 147



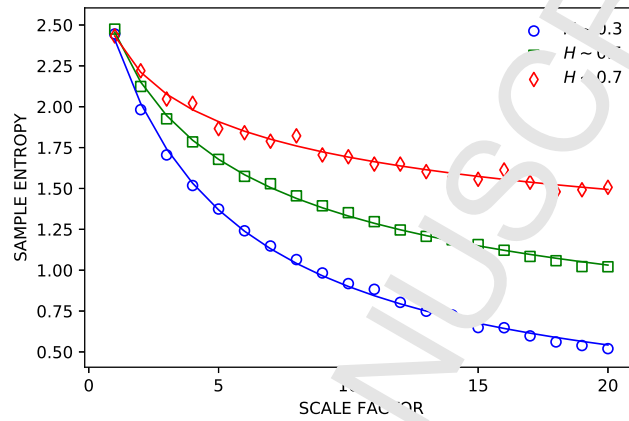


Figure 2: The MSE analysis for three sample time series with the Hurst exponent,  $\sim 0.3$ ,  $\sim 0.5$  and  $\sim 0.7$ . The plot for time series with small Hurst exponent decreases more rapidly than others. The results of fitting with the decreasing  $q$ -exponential are plotted by solid lines. The values of  $q$  parameter for these time series are 2.07, 3.42 and 6.35 respectively.

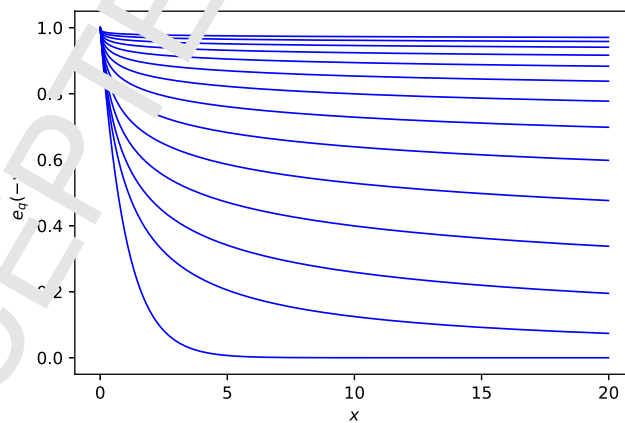


Figure 3: The decreasing  $q$ -exponential function for wide range of  $q$  from 1 to 400. The large value of  $q$  mimics the MSE result for  $1/f$  noise.

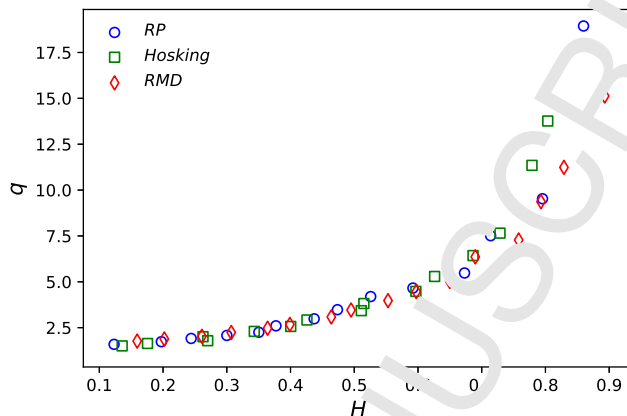


Figure 4: The parameter  $q$  against the Hurst exponent. The value of  $q$  rapidly grows by increasing the Hurst exponent, which ensures us for simulating the behavior of  $1/f$  noise.

148 we put  $m = 2$  and  $r = 0.15\sigma$  where  $\sigma$  is the standard deviation of time series.  
 149 In figure 2 we plot the result for three time series with the Hurst exponent  
 150 equal to,  $\sim 0.3$ ,  $\sim 0.5$  and  $\sim 0.7$ . As is seen in this figure, all plots show  
 151 decreasing behavior but for large value of the Hurst exponent, this behavior  
 152 is more gentle.

153 This behavior could be justified by closer look at the different stages of  
 154 the MSE calculation. In the coarse graining procedure, we eliminate the short  
 155 range correlated (high frequency) noises in time series, hence the variance  
 156 of the resulting coarse grained time series decreases. Since  $r$  is constant,  
 157 decreasing in variance causes that  $n_{m+1}(r, \tau)$  tends to  $n_m(r, \tau)$  and conse-  
 158 quently the sample entropy decreases too [18]. It is expected that the time  
 159 series with lesser Hurst exponent shows more decrease in the variance and  
 160 also in the sample entropy because the high frequency noises are dominated  
 161 for lower values of  $H$ .

162 It is important to mention that the FGN series with  $H = 0.5$  is the same  
 163 as the white noise and as we expected, its result completely overlaps with  
 164 the white noise result.

165 In order to find the relation between the MSE analysis and the Hurst  
 166 exponent, it is required to model the result of analysis with a model func-  
 167 tion. This function should be well fitted to all the FGN results only by  
 168 changing some parameters. The number of these parameters must be at  
 169 minimum. Here we choose, the following function which is defined in terms

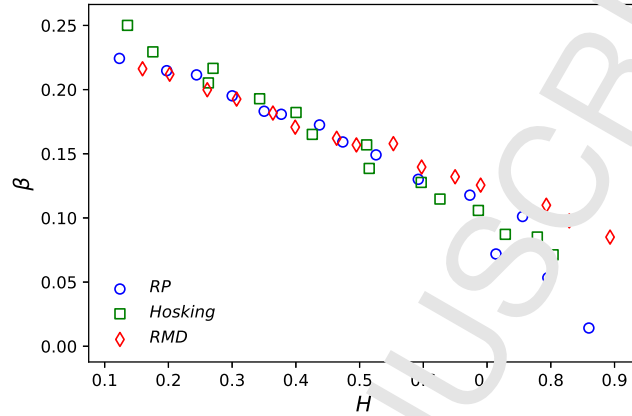


Figure 5: The parameter  $\beta$  against the Hurst exponent. The value of  $\beta$  decreases by increasing the Hurst exponent. It seems that for large value of the Hurst exponent, this decreasing becomes more rapidly.

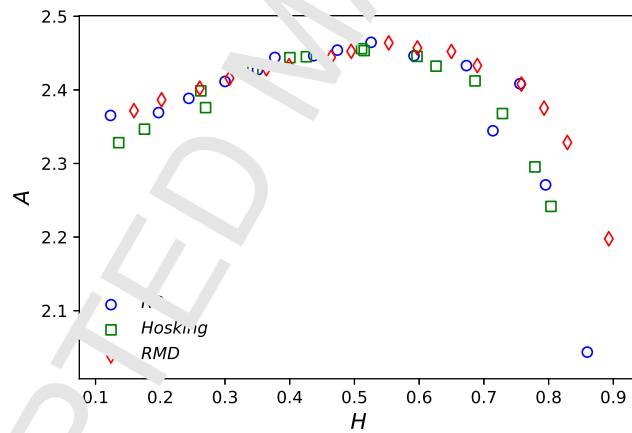


Figure 6: The parameter  $A$  against the Hurst exponent. The behavior of  $A$  is different for  $H < 0.5$  and  $H > 0.5$ . First it increases slowly then rapid decreasing is observed.

170 of  $q$ -exponential function,

$$f(x; A, q, \beta) = Ae_q(-\beta x) = A[1 - \beta(q-1)x]^{\frac{1}{1-q}} \quad (7)$$

171  $A$ ,  $q$  and  $\beta$  are three real positive numbers.  $q$  must be one or greater than  
 172 one. It has decreasing behavior for the model function. In figure 2 we fit  
 173 the results of analysis for the above mentioned three sample time series by  
 174 this model function. As is seen, the analysis results are well fitted to the

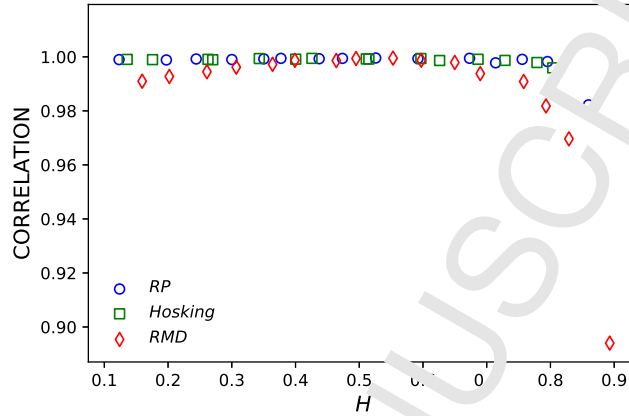


Figure 7: The correlation measure of fitting procedure. The MSE analysis results for time series that are generated by the Hosking algorithm and the Rambaldi-Pinazza method are well fitted by our suggested function.

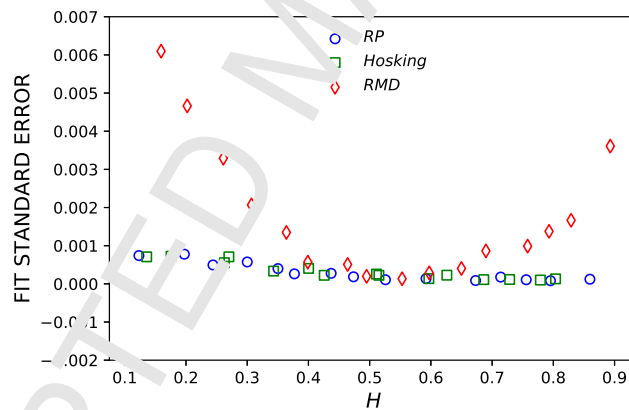


Figure 8: The fit standard error measure of fitting procedure. The MSE analysis results for time series that are generated by the Hosking algorithm and the Rambaldi-Pinazza method are well fitted by our suggested function.

175 suggested function. The values of  $q$  parameter are 2.07, 3.42 and 6.35 for  
 176 time series with the Hurst exponent  $\sim 0.3$ ,  $\sim 0.5$  and  $\sim 0.7$  respectively.  
 177 The larger Hurst exponent leads to the greater value of  $q$  parameter and the  
 178 gentle decrease of the model function.

179 One advantage of using such function, is mimicking behavior of the MSE  
 180 results from white noise to  $1/f$  noise by only one function. This statement is  
 181 illustrated graphically in figure 3. By increasing the value of  $q$ , the function

182 becomes close to the horizontal straight line which is the behavior of the  
 183 obtained result for the  $1/f$  noise.

184 We fit all the results that obtained from the MSE analysis of the FGN time  
 185 series by the mentioned function. Figures 4 and 5 show the fit parameters,  $q$   
 186 and  $\beta$  against the Hurst exponent. We observe increasing behavior for  $q$  and  
 187 decreasing behavior for  $\beta$ . The parameter  $q$  varies between 1.5 for  $H = 0.136$   
 188 to 18.9 for  $H = 0.86$ . Simultaneously,  $\beta$  changes in the range of 0.25 to 0.01.  
 189 The parameter  $A$  softly increases from 2.3 to 2.45 for  $H < 0.5$  then rapid  
 190 decreasing can be observed, see the figure 6. We know that the FGN process,  
 191 when  $H \rightarrow 1$  simulates the  $1/f$  noise. It can be understood through the large  
 192 value of  $q$  and small value of  $\beta$  for time series with large Hurst exponent,  
 193 because the resulting model function becomes very close to the horizontal  
 194 straight line.

195 The fit standard error and correlation are two important measures which  
 196 determine the goodness of fitting. It is worth to note that both of them  
 197 indicate that all the curve fittings are well done. The correlation coefficient  
 198 for nearly all data is more than 0.99 (see figure 7) and the fit standard error  
 199 for most of the fitting procedures are less than 0.001 (see figure 8). The  
 200 results for time series which were generated by the Hosking algorithm and  
 201 the Rambaldi-Pinazza method are better fitted to the model function than  
 202 others.

203 Finally, we mention that for the MSE analysis we use the program written  
 204 by M. Costa *et al.*, which can be downloaded from [physionet.org](http://physionet.org) [26]. In  
 205 fitting procedure we benefit from the curve fitting function in SciPy which is  
 206 the Python library for the scientific and technical computing.

## 207 5. Conclusion

208 The MSE analysis is favorable way for the abnormally detection in phys-  
 209 iological time series. It has been widely used for analysis of the heart beats  
 210 and the electroencephalogram time series. Many diseases were distinguished  
 211 by such analysis without any need to use the invasive methods. This way of  
 212 the analysis of complexity has been applied to technical time series too. In  
 213 this study, we were interested to answer a fundamental question, how MSE  
 214 can be used for detection of the long range dependence in time series. In  
 215 other word, we were looking to explore the relation between the MSE anal-  
 216 ysis result and the Hurst exponent of a time series. We chose the FGN time  
 217 series with different Hurst exponent as our data samples for analysis. It was

218 found that by the function,  $Ae_q(-\beta x)$ , all the results for the FGN time se-  
 219 ries can be modeled. We explored the nice relation between parameters of  
 220 model function;  $q, \beta$  and  $A$ , with the Hurst exponent. We also reported the  
 221 curve fitting measures which demonstrated that the well fitting were done.  
 222 In future, we would interested to find the effect of trends on the result of  
 223 the MSE analysis and apply it on natural time series, particularly the daily  
 224 temperature data and finding a way for the climate classification.

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